

# Subcriticality control elements in a reactor system with an extended plasma source of neutrons with regard for temperature\*

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## Abstract

Materials have been selected for the shim rods and burnable absorbers to compensate for the excessive reactivity of the facility's blanket part and to provide for the possibility of reactivity control in conjunction with a plasma source of neutrons.

Burnable absorber is a layer of zirconium diboride ( $ZrB_2$ ) with a thickness of 100  $\mu m$  applied to the surface of fuel compacts. Boron carbide ( $B_4C$ ) rods installed in the helium flow channels and used to bring the entire system into a state with  $k_{eff} = 0.95$  have been selected as the shim rod material. Throughout its operating cycle, the facility is subcritical and is controlled using the neutron flux from the plasma source.

Verified codes, WIMS-D5B (ENDF/B-VII.0) and MCU5TPU (MCUDB50), as well as a modern system of constants were used for the calculations.

The facility's neutronic performance was simulated with regard for the changes in the inner structure and temperature of the microencapsulated fuel and fuel compact materials caused by long-term irradiation and by the migration of fission fragments and gaseous chemical compounds.

## Keywords

Fusion-fission reactor system, plasma neutron source, criticality, burnable absorber, control and protection system

## Introduction

High-temperature gas-cooled reactors (HTGR), due to the peculiarities of the core design composition and layout, feature advantageous characteristics in terms of nuclear safety and reliability (Alekseyev et al. 2015). In an HTGR, the power unit size is selected such that it would be possible to remove decay heat and to avoid the escape of fis-

sion products from the fuel into the coolant and then into the environment. These systems practically exclude the core melting, especially for cores composed of ceramics and graphite which are capable to withstand continuously temperatures in a range of 1000 to 3000 K (Alekseyev et al. 2005). No loss of coolant leads to an abrupt temperature growth which is explained by the core's high heating power. Using helium as the coolant excludes chemical

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reactions with the fuel and structural materials, and does not lead to problems involving phase transformations (Grebennik et al. 2008). Besides, HTGRs have a high negative temperature reactivity coefficient (Alekseyev et al. 2005, 2015, Grebennik et al. 2008, Shamanin 2008), this ensuring the reliability of operation during transients and in power maneuvering modes.

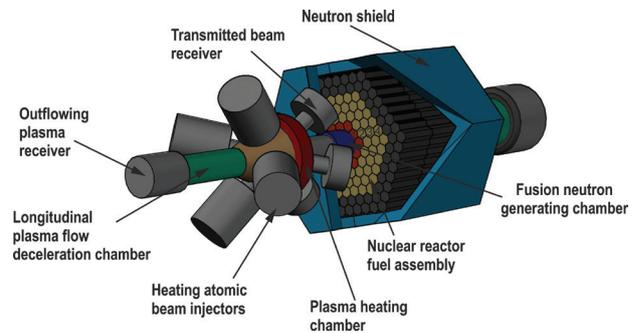
Conventional reactor systems operate in a critical mode with the fission chain reaction controlled by influencing the balance of neutrons in the core's breeding environment. Operation of the facility under consideration, which consists of a blanket (the facility's energy generating part) and an external plasma source of fusion neutrons, differs from the traditional operating modes of the reactor, still the principle of control is basically the same and consists in influencing the balance of neutrons in the facility's active blanket part. The effective multiplication factor for the loaded blanket of the facility under consideration is made up for to the values 0.95 to 0.98 and is maintained at the preset level using criticality control elements and systems. The plasma source generates extra neutrons and ensures, in conjunction with the control system, safe operation of the entire facility (Arzhannikov et al. 2019a, Prikhodko and Arzhannikov 2020), improving so nuclear safety since, with this source being out of the generation mode, the whole of the blanket part will render deeply subcritical ( $k_{\text{eff}} < 0.95$ ).

The efficiency of compensating systems depends on a number of parameters investigated by the authors in (Kalin et al. 2008, Rusinkevich 2016, Bedenko et al. 2019) and in other papers. Investigating the influence the changing thermophysical properties of the facility's blanket part have on its neutronic performance and the efficiency of the compensating system (Rusinkevich 2016, Bedenko et al. 2019) is of the utmost interest. Given these factors and the operating peculiarities of the subcritical facility under investigation, a vital objective discussed in this paper is to select the burnable absorber materials and to develop the shim rod (SR) model and layouts in conditions of long-term operation and high burn-up.

## Materials and methods

Fig. 1 shows the layout of the facility under consideration (Arzhannikov et al. 2019a). Its active part's region adjoining the axis is substituted by a cylindrical vacuum chamber containing high-temperature plasma that generates extra neutrons as a result of D-D and D-T thermonuclear reactions. This plasma part has a chamber mated to it for injecting beams of high-energy neutral deuterium and tritium atoms (Anikeev et al. 2015, Yurov et al. 2016). The magnetic field in these two portions of the vacuum chamber, which contain high-temperature plasma, thermally isolates the plasma from the chamber walls in the radial direction (Arzhannikov et al. 2019). The parameters of the magnetic field and the plasma source of neutrons for making up the facility's blanket part, as well as the blanket configuration and composition are presented in (Arzhannikov et al. 2019, 2019a, 2020).

Structurally, the blanket consists of hexagonal graphite blocks having channels for the accommodation of fuel and injection (flow) of helium surrounded by two rows of graphite blocks which perform the reflector function. At the top and at the bottom, the blanket part is shielded with graphite blocks installed in one row. The fuel block has 76 small-diameter channels for fuel and seven large-diameter channels for helium. The width across flats is 0.207 m, and the height is 0.8 m. The graphite blocks that shield the blanket part at the top and at the bottom have a width across flats of 0.207 and a height is 0.3 m. The microencapsulated (coated particle) fuel for the graphite fuel compacts represents spherical kernels of  $(\text{Th,Pu})\text{O}_2$  with a diameter of  $0.350 \times 10^{-3}$  m covered with PyC and  $\text{Ti}_3\text{SiC}_2$  layers of the thickness  $0.90 \times 10^{-4}$  m and  $0.35 \times 10^{-4}$  m respectively. The coated particle fuel is dispersed into the graphite matrix of the fuel compact to the surface of which an outer force coat of SiC is additionally applied. The fuel compact diameter is  $10.17 \times 10^{-3}$  m, the height is  $20.10^{-3}$  m, and the thickness of the outer SiC layer is  $0.3 \times 10^{-3}$  m.



**Figure 1.** A subcritical facility with a plasma neutron source (Arzhannikov et al. 2019a).

## Calculation of the temperature in the fuel compact

The BOL maximum fuel temperature was estimated based on a condition that there was bulk heat generation,  $q_v(z,r)$ , and surface heat generation,  $q_s(z,r)$ , in the fuel block. With the rated power of the facility being 60 MW, the maximum values  $q_v^{\text{max}}(z,r)$  and  $q_s^{\text{max}}(z,r)$  are respectively equal to  $8.41 \times 10^4$  and  $2.01 \times 10^4$  kW/m<sup>3</sup>.

The computational studies presented in (Arzhannikov et al. 2020) show that the facility's fuel life, with a thermal power of 60 MW, is 300 eff. days. Such long-term operation leads to a major change in the coated particle fuel's nuclide composition and structure, the accumulation and migration of fission fragments and gaseous chemical compounds, and, as a consequence, to a change in the fuel's thermophysical properties (Bedenko et al. 2019) and temperature (Rusinkevich 2016).

A code, MCU5TPU (MCUDB50), is used to calculate the evolution of the fuel's nuclide composition and to estimate the radiation dose and the neutron fluence. The MCU-5 geometrical module makes it possible to simula-

te 3D systems with a geometry of any complexity when using a combined approach based on describing complex systems by combinations of elementary bodies and surfaces (MCU Project). The MCU-5 nuclide library includes an extensive list of isotopes and allows calculating the evolution of the nuclide composition and the criticality.

The effects of the fuel component's nuclide composition and the blanket part's thermophysical properties and temperature on the neutronic performance and efficiency of the compensating system have been calculated for an equivalent 2D cylindrical cell based on models in (Kalin et al. 2008, Hales et al. 2013, Degaltsov et al. 1987). These models make it possible to take additionally into account the mechanisms of the fission fragment diffusion and the migration of gaseous chemical compounds formed in the interaction of the oxygen released in coated particle fuel with fission fragments and the initial coat material. The migration and diffusion of fission products and the chemical compounds formed into the fuel compact matrix have been calculated with regard for the retarding capacity of the fuel kernel's principal diffusion barrier (Rusinkevich 2016, Ponomarev-Stepnoy et al. 2009).

## Calculation of shim elements

The blanket part of the facility is a modified HTGR core with a neutron spectrum close to the epithermal spectrum. Compounds employed traditionally in high-temperature reactors ( $B_4C$ ,  $B_4C-SiC$ ,  $Dy_2O_3TiO_2$ ,  $CrB_2+Al$ ,  $Gd_2O_3$ ,  $ZrB_2$ ,  $AgInCd$ ,  $Mo+Eu_2O_3$ ,  $Hf-Zr$ , and others) were therefore used to select the effective materials for the shim rods (SR) and the burnable absorbers (BA) (Kalin et al. 2008).

Passive reactivity should be achieved with the use of SRs and BAs in conjunction with a plasma neutron generator which, together with the SR and BA system, is required to compensate for the effects caused by the nuclear fuel burn-up, and by the blanket slagging and poisoning in the process of the startup and during long-term operation. This was achieved in a 69-group diffusion approximation by the joint use of the WIMS-D5B (ENDF/B-VII.0) code (Pazirandeh et al. 2011) and an iterative method used to solve the neutron transport equation (Shamanin et al. 2017). The calculations take into account the effects leading to a change in the thermophysical properties of the blanket and in the temperature of the compact's fuel part.

After identifying the neutron absorbing material, which is effective in the operating neutron spectrum, the SRs and BAs are calculated. The best possible SR arrangement in the facility's blanket is also generated.

## Results

### Temperature calculation

The calculation, the results of which are presented in Figs 2, 3, has been performed for the facility's most heated fuel block as of the BOL. The outlet coolant tempera-

ture is limited to a value of 1300 K, and the permissible helium flow rate is 50 m/s. The temperature choice was not random since it is investigated if it is additionally possible to use high-temperature heat for implementing the technology of methane conversion to hydrogen.

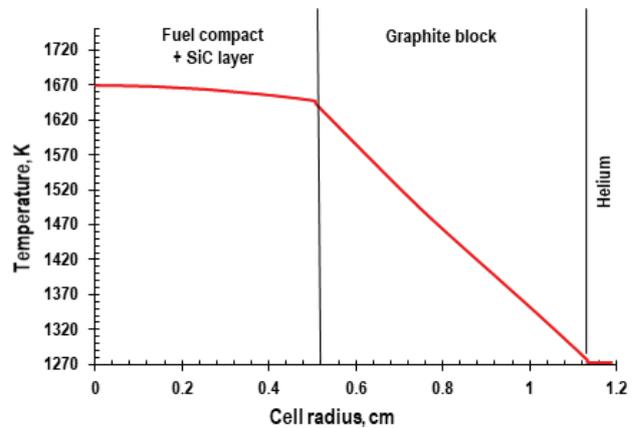


Figure 2. Temperature distribution by calculation cells.

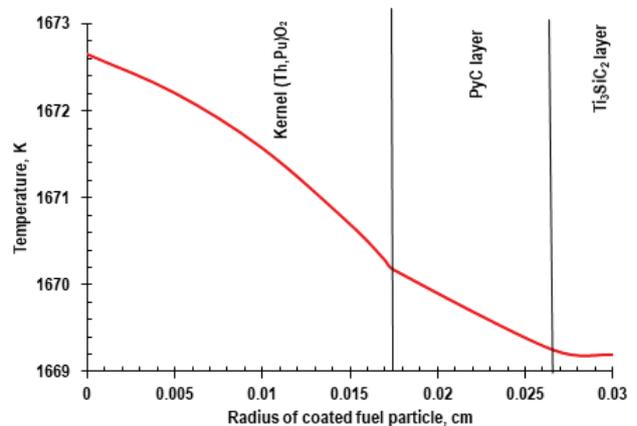


Figure 3. Temperature distribution in microencapsulated fuel.

The accumulation of fission fragments and gaseous compounds and the migration of these in conditions of long-term irradiation lead to a temperature growth in the kernel (Rusinkevich 2016, Hales et al. 2013, Degaltsov et al. 1987, Ponomarev-Stepnoy et al. 2009). As shown by the calculation results, the temperature in the kernels will grow by 40 K 1500 days after the facility startup. By the end of the fuel life, the temperature in the kernels will cumulatively increase by 70 K. Therefore, the SR and BA materials need to be selected with regard for structural and other changes leading to a temperature growth and affecting the efficiency of the compensating system and the blanket's neutronic parameters.

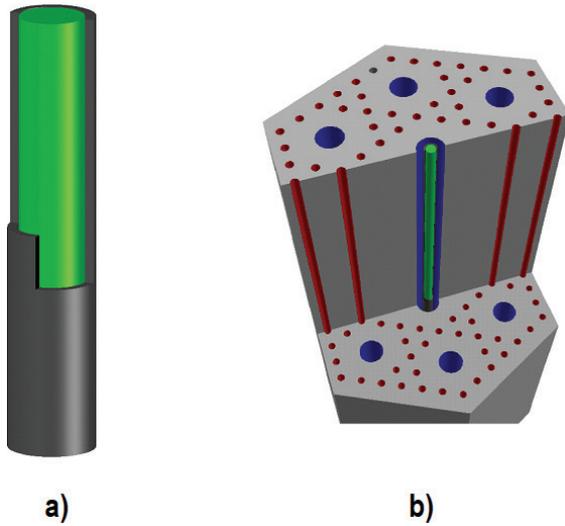
The results of investigating the coated particle fuel and fuel compact materials have shown that the conditions of operation do not exceed the permitted working maximums (Degaltsov et al. 1987, Miller 1995) with the fast neutron fluence being  $\sim 1 \times 10^{18} \text{ m}^{-2}$ . In these conditions, the radiation changes in the kernel and in the coats are minor, and the strength characteristics of the main diffusion barriers are stable (Nappé et al. 2011, Katoh Yutai 2014). The additional safety barrier is formed by the graphite matrix and by the outer force coat.

## Compensating materials and arrangement of compensating elements

An analysis of the findings has shown that the best reactivity compensation options are those with a  $ZrB_2$  coat of the thickness  $0.1 \cdot 10^{-3}$  m applied to the surface of the fuel compacts and with  $B_4C$  (Table 1) used as the SR material (Fig. 4). Table 1 presents the results of calculating the weight of the SR installed at the facility's center expressed in relative units. For comparison,  $Gd_2O_3$  and  $Eu_2O_3$  SRs of 0.01, 0.03 and 0.05 m in diameter have been calculated.

**Table 1.** Results of calculating the weight of one CPS rod

| Material            | $k_{eff}$ , initial | $k_{eff}$ , with rod | $\rho$ of system | $d\rho$ (weight of one rod) |
|---------------------|---------------------|----------------------|------------------|-----------------------------|
| Rod diameter 0.01 m |                     |                      |                  |                             |
| $B_4C$              | 1.207558            | 1.207012             | 0.171508         | 0.000375                    |
| $Eu_2O_3$           | 1.207558            | 1.207119             | 0.171581         | 0.000301                    |
| $Gd_2O_3$           | 1.207558            | 1.207235             | 0.171661         | 0.000222                    |
| Rod diameter 0.03 m |                     |                      |                  |                             |
| $B_4C$              | 1.207558            | 1.205919             | 0.170757         | 0.001126                    |
| $Eu_2O_3$           | 1.207558            | 1.206108             | 0.170887         | 0.000996                    |
| $Gd_2O_3$           | 1.207558            | 1.206296             | 0.171016         | 0.000866                    |
| Rod diameter 0.05 m |                     |                      |                  |                             |
| $B_4C$              | 1.207558            | 1.205016             | 0.1701355        | 0.001747                    |
| $Eu_2O_3$           | 1.207558            | 1.20537              | 0.1703792        | 0.001503                    |
| $Gd_2O_3$           | 1.207558            | 1.20542              | 0.1704136        | 0.001469                    |



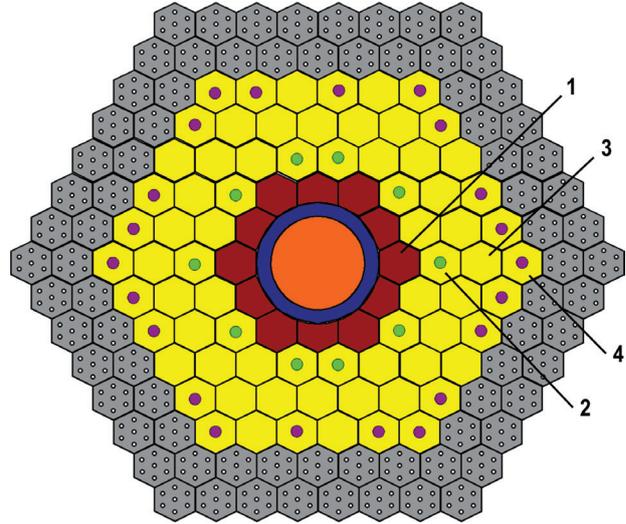
**Figure 4.** Model of the channel layout with a movable SR: a) – SR channel; b) – channel arrangement in the fuel block.

Using  $ZrB_2$  makes it possible to reduce the effective multiplication factor to a value of  $k_{eff} = 1.069833$ . Further blanket subcriticality of up to  $k_{eff} = 0.95$  was achieved using SRs of 0.01 m in diameter accommodated in the helium channels. The required number of rods is 138. The SR arrangement in the fuel block is presented in Fig. 4.

The SR channel material is steel of the KhN55M-VTs-VI(ID) (ChS57-VI(ID)) grade (Heat-resistant nickel-based alloy) which has proved itself to perform well in conditions of long-term in-pile irradiation.

Simulation of various options for the SR arrangement in blocks and for the block positions in the blanket made it possible to obtain the optimum SR arrangement

presented in Fig. 5 (the figures show the fuel block row numbers). Therefore, the blanket's row 2 contains 10 blocks with fully loaded SRs, and row 4 contains 22 such blocks. There are 5 SRs installed in each block in the blanket's row 2, and there are 4 SRs installed in each block in row 4. Such arrangement leads to the effective multiplication factor reduced to approximately 0.9494, which corresponds to the required value for the facility operation.



**Figure 5.** Arrangement of blocks with SRs in the facility's active part.

## Discussion

The arrangement of fuel blocks with SRs shown in Fig. 5 allows reducing the material costs (the weight of the BAs for the SRs out of those considered in Table 1 is as small as possible), leveling off the power profile, and achieving the required subcriticality.

We shall note that the use of a plasma neutron source improves the facility's nuclear safety since switching off the injection of neutral atoms causes the generation of neutrons to drop by approximately a half for the initial 2.5 ms and by a factor of 20 more for the further 5 ms (Prihodko and Arzhannikov 2020), that is, much faster than in the core of a traditional reactor.

## Conclusion

The most advantageous method for reactivity compensation is to apply a  $ZrB_2$  coat of the thickness  $0.1 \cdot 10^{-3}$  m to the surface of the fuel compact and to use  $B_4C$  SRs of 0.01 m in diameter installed in the coolant channels. Coat application and sintered  $B_4C$  manufacturing technologies are successfully developed by research teams at the Tomsk Polytechnic University.

The burnable absorber and shim rod materials have been chosen with regard for the changes in the inner structure of the microencapsulated fuel and the fuel com-

pact caused by long-term irradiation and by the migration of fission fragments and gaseous chemical compounds affecting the thermophysical properties and the temperature of the compact's fuel part.

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## References

- Alekseyev PN, Kukharkin NYe, Udyanskiy YuN, Schepetina TD, Belov IA, Subbotin SA, Sedov AA, Dudnikov AA, Kashka MM, Bashayev VV (2005) Advanced nuclear power plants with fuels based on coated fuel particles for nuclear ships. *Atomnaya energiya* 99(1): 3–8. <https://doi.org/10.1007/s10512-005-0229-z> [in Russian]
- Alekseyev PN, Subbotin SA, Stukalov VA, Shchepetina TD (2015) A system of small nuclear power plants as a factor of national security. *Akademiya energetiki* 2(64): 74–79. [in Russian]
- Anikeev AV, Bagryansky PA, Beklemishev AD, Ivanov AA, Korobeinikova OA, Kovalenko YuV, Lizunov AA, Maximov VV, Murakhtin SV, Pinzhenin EI, Prikhodko VV, Savkin VYa, Soldatkina EI, Solomakhin AL, Yakovlev DV, Zaytsev KV (2015) The GDT experiment: status and recent progress in plasma parameters. *Fusion Science and Technology* 68(1): 1–7. <https://doi.org/10.13182/FST14-867>
- Arzhannikov A, Bedenko S, Shmakov V, Knyshev V, Lutsik I, Prikhodko V, Shamanin I (2019) Gas-cooled thorium reactor at various fuel loadings and its modification by a plasma source of extra neutrons. *Nuclear Science and Techniques* 30(12): 1–11. <https://doi.org/10.1007/s41365-019-0707-y>
- Arzhannikov AV, Shamanin IV, Bedenko SV, Prikhodko VV, Sinitsky SL, Shmakov VM, Knyshev VV, Lutsyk IO (2019a) A hybrid thorium reactor plant with a magnetic trap source of fusion neutrons. *Izvestiya vuzov. Yadernaya energetika* 2: 43–54. <https://doi.org/10.26583/npe.2019.2.04> [in Russian]
- Arzhannikov AV, Shmakov VM, Modestov DG, Bedenko SV, Prikhodko VV, Lutsik IO, Shamanin IV (2020) Facility to study neutronic properties of a hybrid thorium reactor with a source of thermonuclear neutrons based on a magnetic trap. *Nuclear Engineering and Technology* 52(11): 2460–2470. <https://doi.org/10.1016/j.net.2020.05.003>
- Bedenko S, Karengin A, Ghal-Eh N, Alekseev N, Knyshev V, Shamanin I (2019) Thermo-Physical Properties of Dispersion Nuclear Fuel for a New-Generation Reactors: A Computational Approach. *AIP Conference Proceedings* 2101(1). <https://doi.org/10.1063/1.5099594>
- Degaltsov YuG, Ponomarev-Stepnoy NN, Kuznetsov VF (1987) Behavior of High-temperature Nuclear Fuel during Irradiation. Moscow. Energoatomizdat Publ., 208 pp. [in Russian]
- Grebennik VN, Kukharkin NYe, Ponomarev-Stepnoy NN (2008) High-temperature Gas-cooled Reactors as an Innovative Field of Nuclear Power Development. Moscow. Energoatomizdat Publ., 136 pp. [in Russian]
- Hales JD, Williamson RL, Novascone SR, Perez DM, Spencer BW, Pastore G (2013) Multidimensional Multiphysics Simulation of TRISO Particle Fuel. *Journal of Nuclear Materials* 443: 531–543. <https://doi.org/10.1016/j.jnucmat.2013.07.070>
- Heat-resistant nickel-based alloy, grade HN55MVTs(VI) (2020) Heat-resistant nickel-based alloy, grade HN55MVTs(VI). <http://www.crisp-prometey.ru/science/steel/heat-resistant-alloy-nickel-based-HN55MVTs-VI-for-high-power-plants-withgas-cooled.aspx>. [accessed Aug. 05, 2020] [in Russian]
- Kalin BA, Platonov PA, Chernov II, Shtrombakh YaI (2008) *Physical Materials Science. Vol. 6. Part 2. Nuclear Fuel Materials*. Ed. by B.A. Kalin. Moscow. MIFI Publ., 604 pp. [in Russian]
- Katoh Y, Snead LL, Cheng T, Shih C, Daniel Lewis W, Koyanagi T, Hinoki T, Henager Jr CH, Ferraris M (2014) Radiation-tolerant joining technologies for silicon carbide ceramics and composites. *Journal of Nuclear Materials* 448(1–3): 497–511. <https://doi.org/10.1016/j.jnucmat.2013.10.002>
- MCU Project (2020) MCU Project. Monte Carlo Simulation of Particle Transport Process. <https://mcuproject.ru/rabout.html> [accessed Aug. 05, 2020] [in Russian]
- Miller GK (1995) Stresses in a spherical pressure vessel undergoing creep and dimensional changes. *International Journal of Solids and Structures* 32(14): 2077–2093. [https://doi.org/10.1016/0020-7683\(94\)00197-5](https://doi.org/10.1016/0020-7683(94)00197-5)
- Nappé JC, Monnet I, Grosseau Ph, Audubert F, Guilhot B, Beauvy M, Benabdesselam M, Thomé L (2011) Structural changes induced by heavy ion irradiation in titanium silicon carbide. *Journal of Nuclear Materials* 409(1): 53–61. <https://doi.org/10.1016/j.jnucmat.2010.12.235>
- Pazirandeh A, Ghaseminejad S, Ghaseminejad M (2011) Effects of various spacer grid modeling on the neutronic parameters of the VVER-1000 reactor. *Annals of Nuclear Energy* 38: 1978–1986. <https://doi.org/10.1016/j.anucene.2011.04.020>
- Ponomarev-Stepnoy NN, Makarov VM, Ivanov AS, Belov IA, Rusinkevich AA, Lindemer T, McEachern D, Razvi J (2009) Evaluation of the thermodynamics of deep burn-up HTGR fuel with plutonium kernels. *Proc. of the 6<sup>th</sup> International Topical Meeting on High Temperature Reactor Technology, HTR 2008*(1): 257–262. <https://doi.org/10.1115/HTR2008-58081>
- Prikhodko VV, Arzhannikov AV (2020) Simulations of fusion neutron source based on the axially symmetric mirror trap for the thorium hybrid reactor. *Journal of Physics: Conference Series* 1647. <https://doi.org/10.1088/1742-6596/1647/1/012004>
- Rusinkevich AA (2016) Thermodynamic Effects in the Transfer of Fission Products in Coated Particle Fuel with High Burn-up Values. *Cand. tech. sci. diss. Moscow. NRC Kurchatov Institute Publ.*, 135 pp. [in Russian]
- Shamanin IV (2008) VTGR reactors with thorium-bearing nuclear fuel: neutronic advantages // *Alternativnaya energetika i ekologiya*. *ISJAE* 11: 48–52. [in Russian]
- Shamanin IV, Bedenko SV, Nesterov VN, Lutsik IO, Prets AA (2017) Solution of a neutron transport multigroup equation system for subcritical systems. *Izvestiya vuzov. Yadernaya energetika* 4: 38–49. <https://doi.org/10.26583/npe.2017.4.04> [in Russian]
- Yurov DV, Prikhodko VV, Tsidulko YuA (2016) Nonstationary model of an axisymmetric mirror trap with nonequilibrium plasma. *Plasma Physics Reports* 42(3): 210–225. <https://doi.org/10.1134/S1063780X16030090>